



Monte Carlo Simulations of Relativistic Runaway Electrons and Terrestrial Gamma Ray Flashes (TGF)

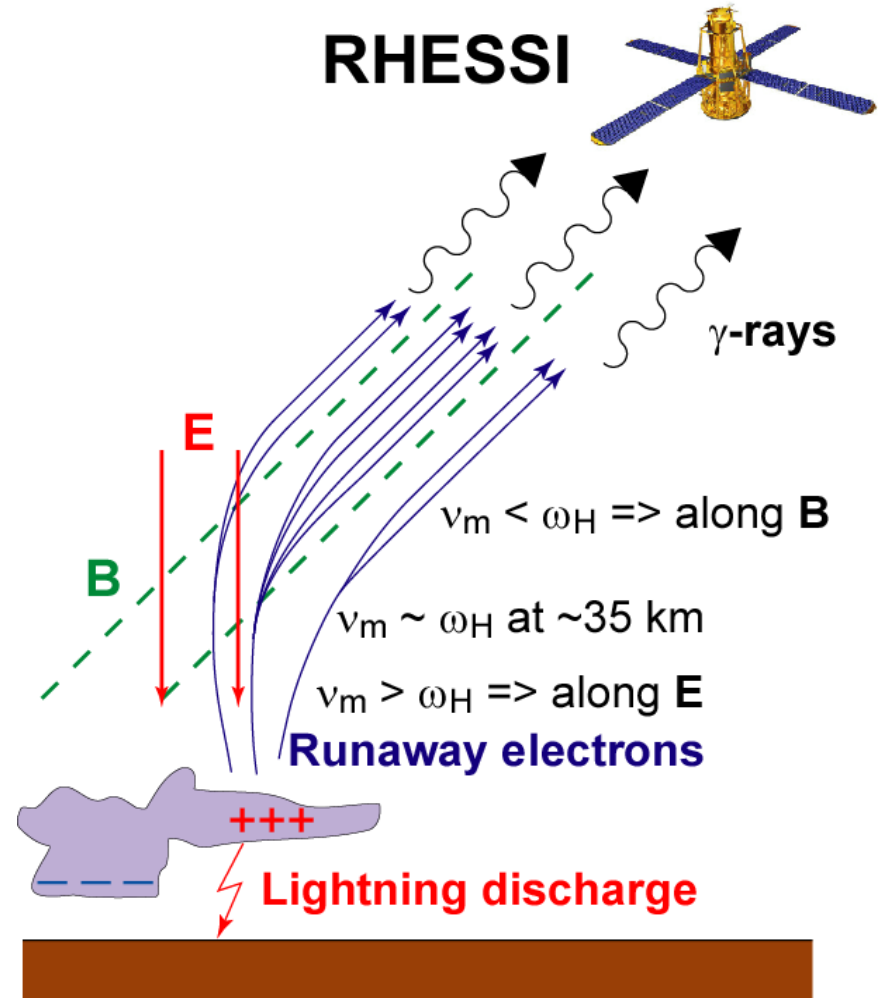
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Terrestrial Gamma Flashes (TGF)

- **~1 ms duration**
- **Energies of up to 20 MeV**
- **Hard spectrum**
- **First observed by BATSE, ~1/month**
- **RHESSI detects ~10-20/month, => ~50/day globally**

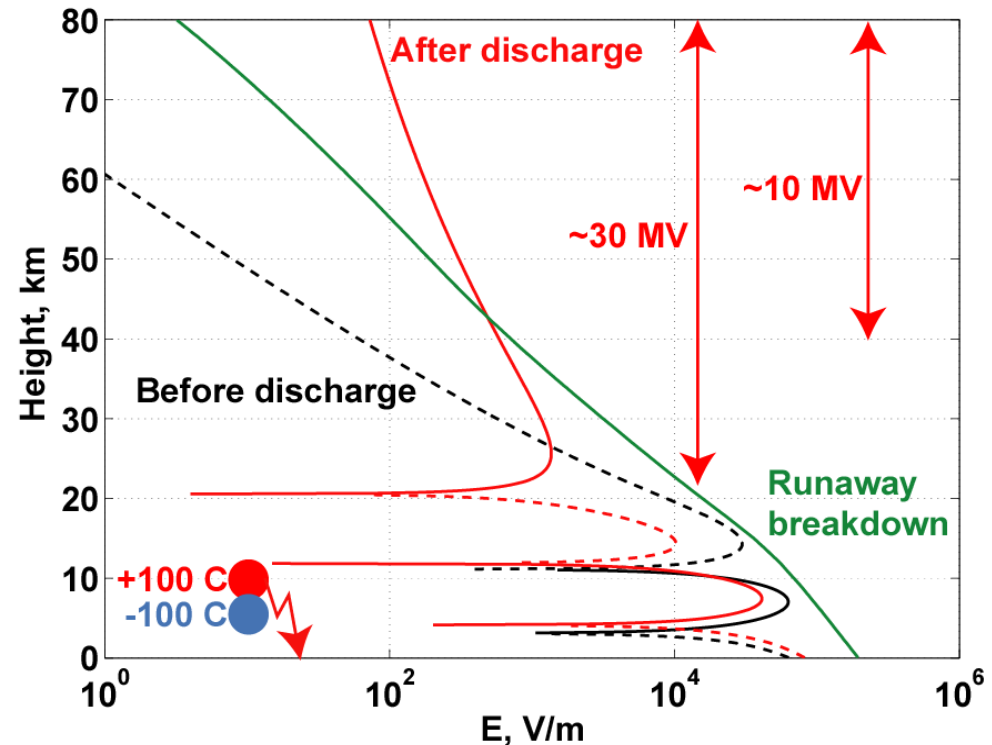
- Electrons are accelerated by E field following a +CG discharge
- Number of electrons is increased due to ionization
- Gamma rays are produced by bremsstrahlung from energetic electrons (~35 MeV energy)
- Emission is forward-directed
- Direction of electron motion is determined by elastic collision rate and B.





Post-discharge electric field

- Exceeds relativistic runaway threshold
- Accelerates electrons upward
- In atmosphere with exponential conductivity profile: 📄

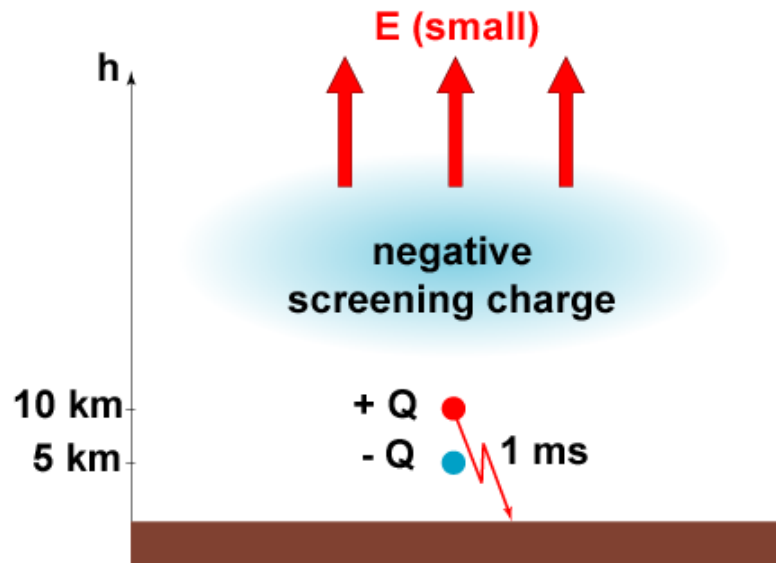


$$\Phi_{\text{disch}} = \frac{Q}{4\pi\epsilon_0} \frac{\exp[-(r+z)/2H] - 1}{r}$$

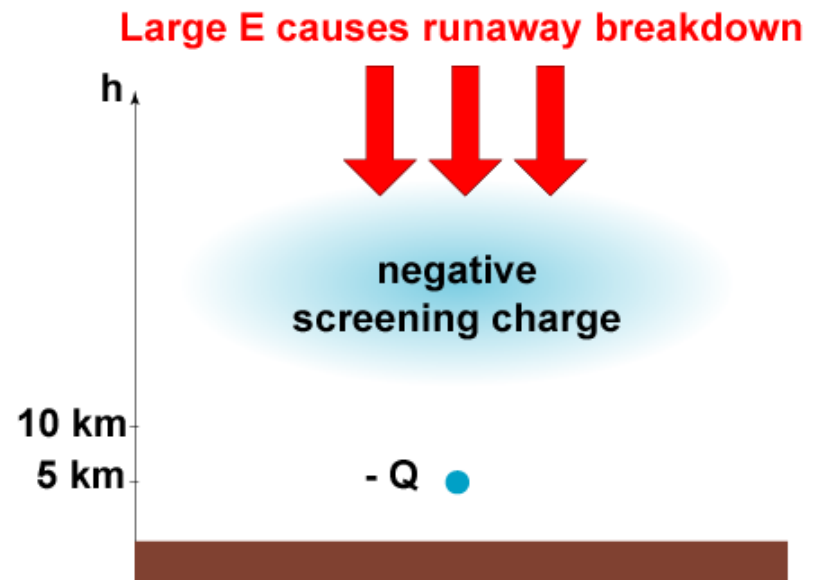


Quasi-Electrostatic Field

BEFORE DISCHARGE

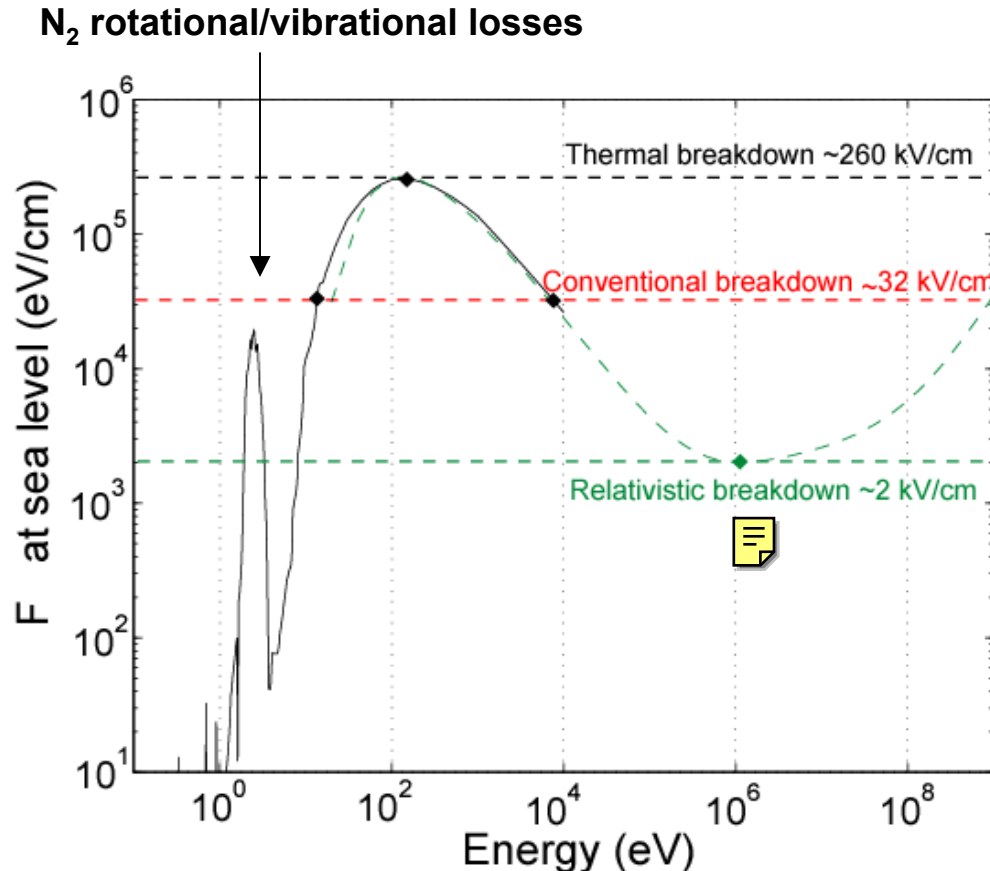


AFTER DISCHARGE





Relativistic Runaway Mechanism and Dynamic Friction Force



$$F = \sum N \sigma_i(v) \Delta \varepsilon_i$$

43 inelastic processes:

- Rotational, vibrational, electronic level excitations
- Dissociative losses
- Ionization

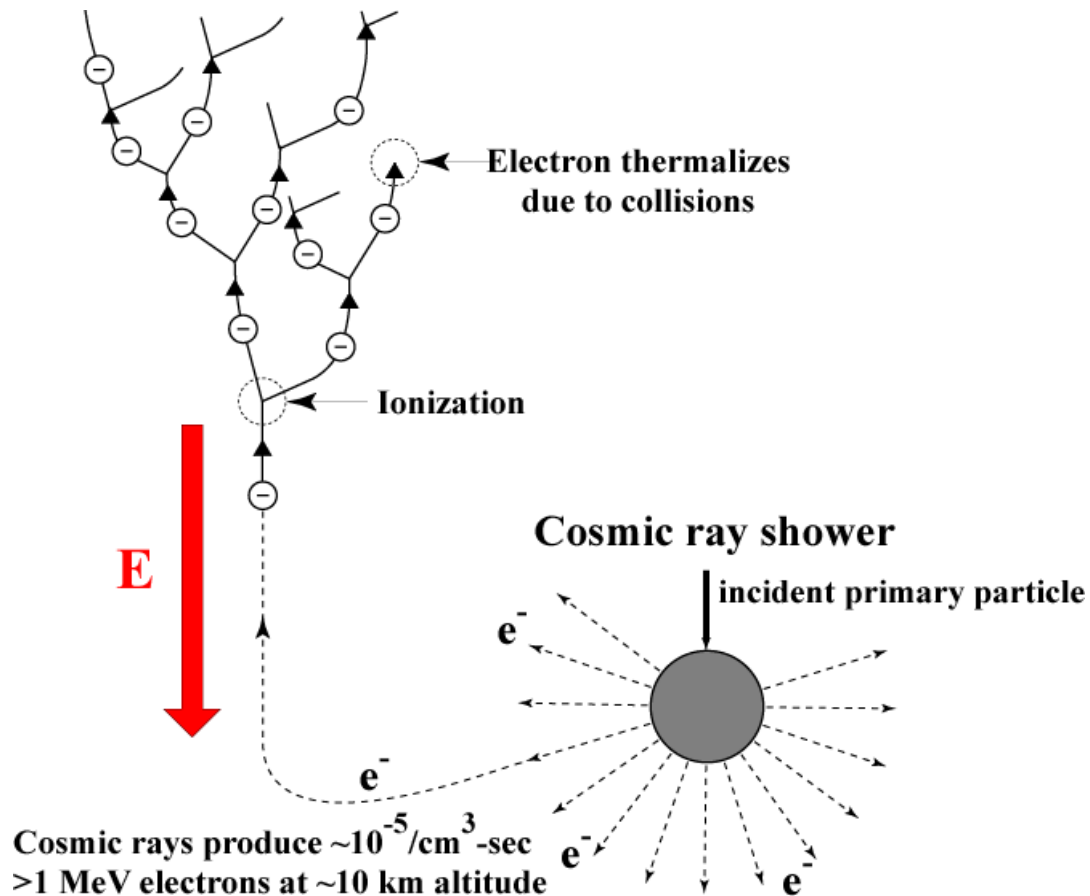
In townsend (1 Td = 10^{-21} V-m²):

$E/N = 8$ Td, 130 Td, 1000 Td



Relativistic Runaway Electron Avalanche

- Cosmic ray primaries
- Ionization with production of relativistic electrons
- Acceleration





Monte Carlo simulations

Method:

- Relativistic motion

$$\frac{d\mathbf{p}}{dt} = -eE - \frac{e}{m\gamma} \mathbf{p} \times \mathbf{B} + \mathbf{G}(t)$$

- $\mathbf{G}(t)$ includes inelastic energy losses and elastic scattering (but excludes energy losses from ionization)
- New electrons from ionization

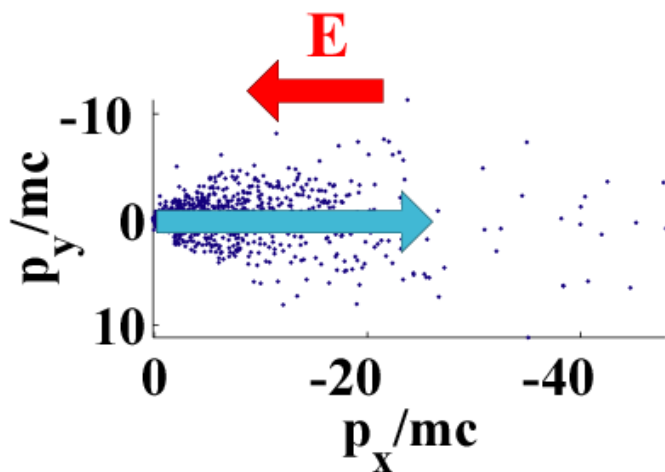
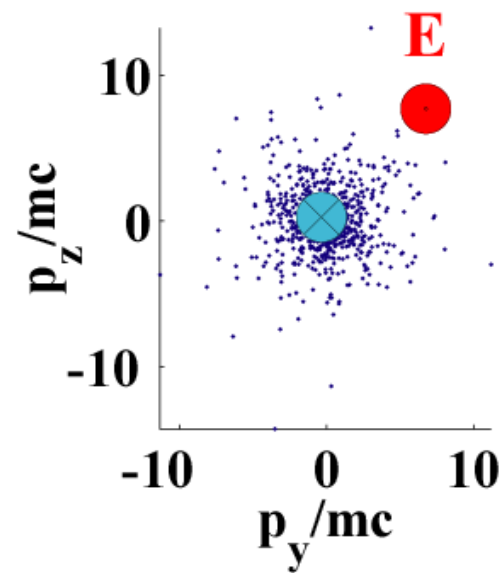
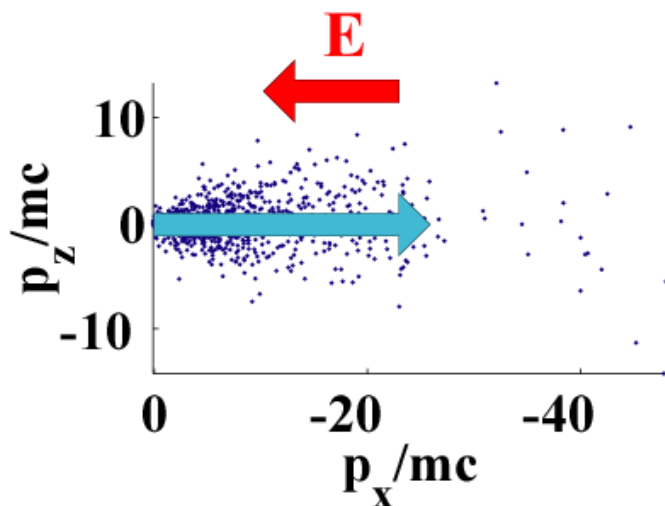
Results:

- Electron distribution
- Avalanche growth rate
- Drift velocity



Electrons in momentum space

$$\mathbf{E} = 5 \mathbf{E}_0 \hat{\mathbf{e}}_x, \mathbf{B} = 0$$



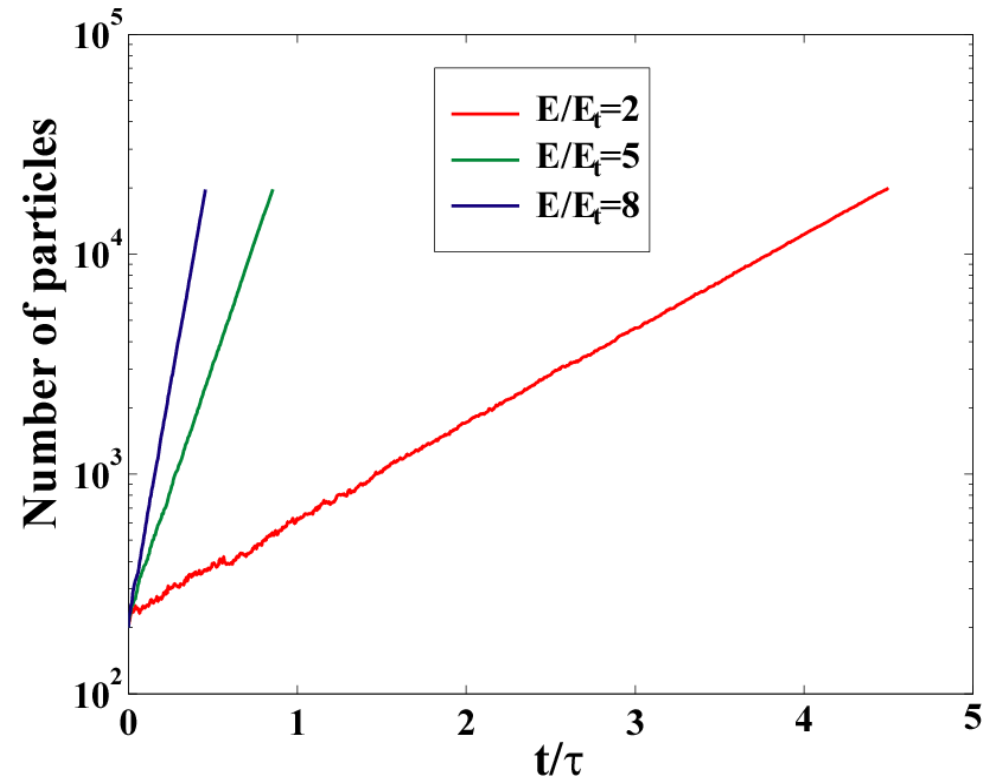


Growth rate

- $N = N_0 e^{Rt}$
- $R = R_0 / \tau$
- Proportional to atmosphere density N_m

$$\tau = (2\pi N_m Z_m r_0^2 c)^{-1}$$

- r_0 – classical electron radius
- Z_m – molecule charge
- $R_0 \sim (\delta - 1) + 0.04(\delta - 1)^2$,
 $\delta = E/E_t$





Fluid modelling of runaway avalanche above thundercloud

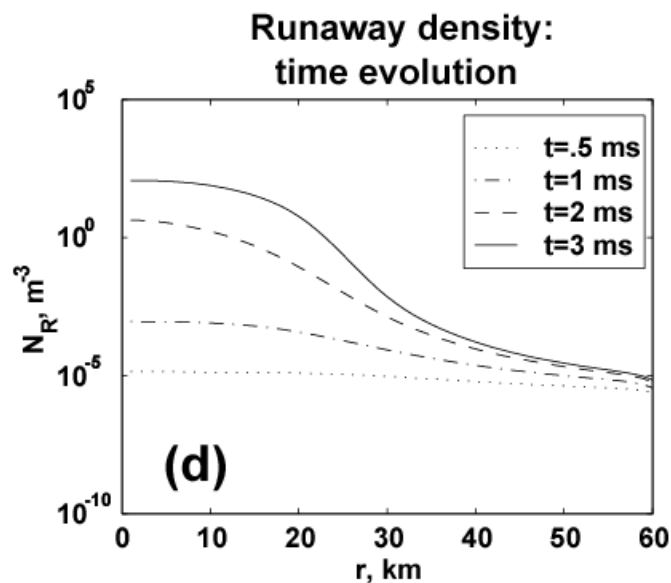
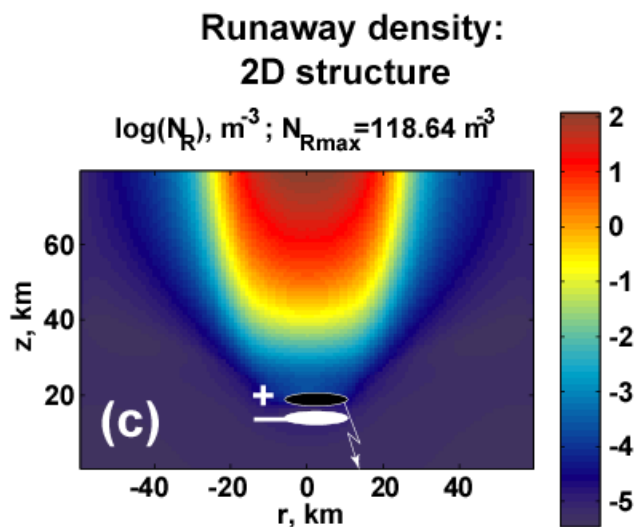
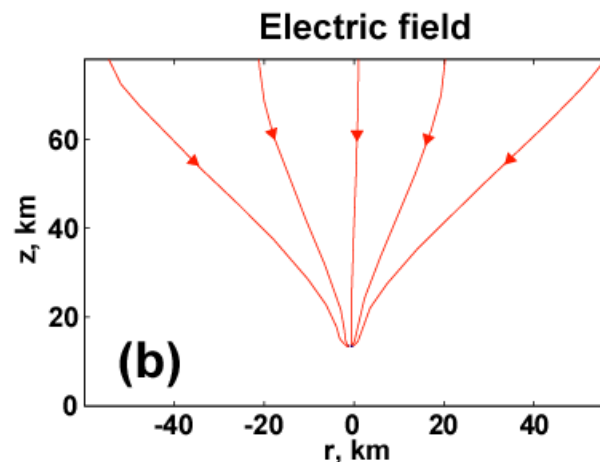
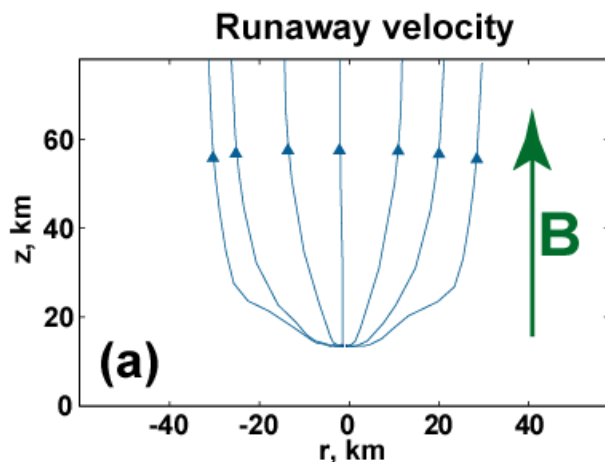
$$\frac{\partial N}{\partial t} + \nabla(vN) = RN + S$$

- **Notations:**
 - v – drift velocity
 - $R(E/E_t) \sim N_m$ – avalanche rate
 - $S \sim N_m$ – source from cosmic rays, $= 10 \text{ m}^{-3} \text{ s}^{-1}$ at 10 km
- **Cylindrical**
 - Vertical magnetic field
- **Cartesian**
 - Arbitrary direction of magnetic field
 - Horizontally extended thundercloud



Fluid modelling results: Cylindrical

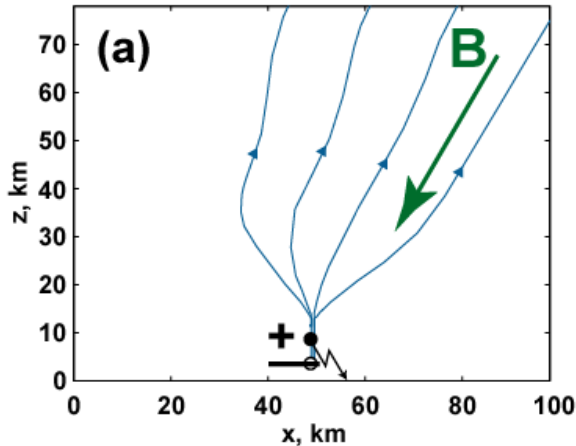
$Q=420$ C, $h_f=20$ km, $t=3$ ms



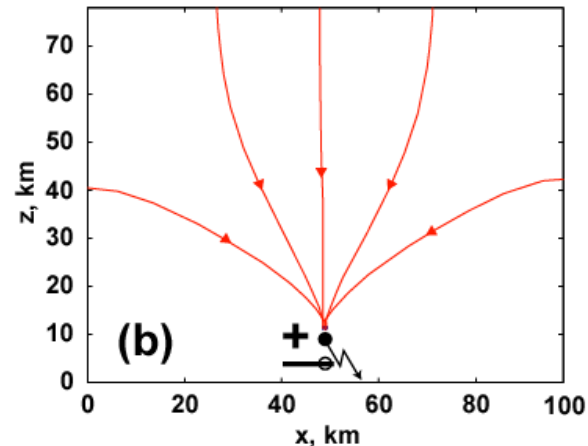
Fluid modelling results: Cartesian

$Q=12 \text{ C/km}$, $h_f=10 \text{ km}$, $t=3 \text{ ms}$

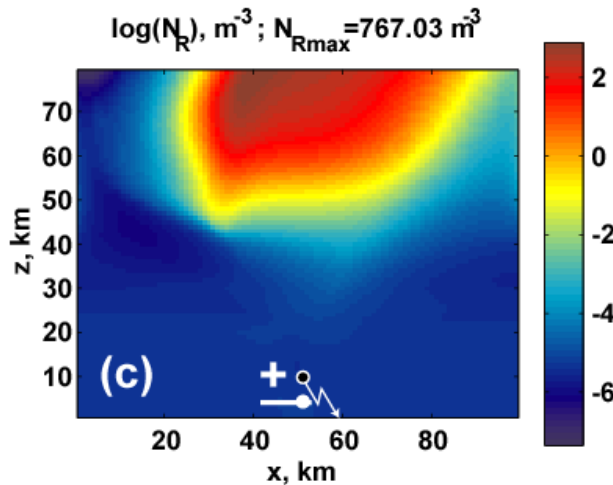
Runaway velocity



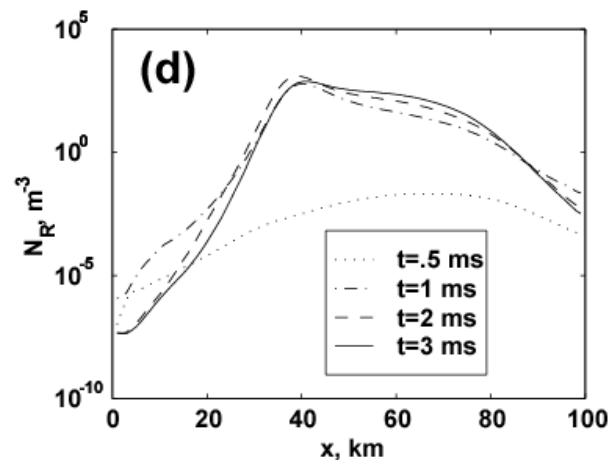
Electric field



Runaway density:
2D structure

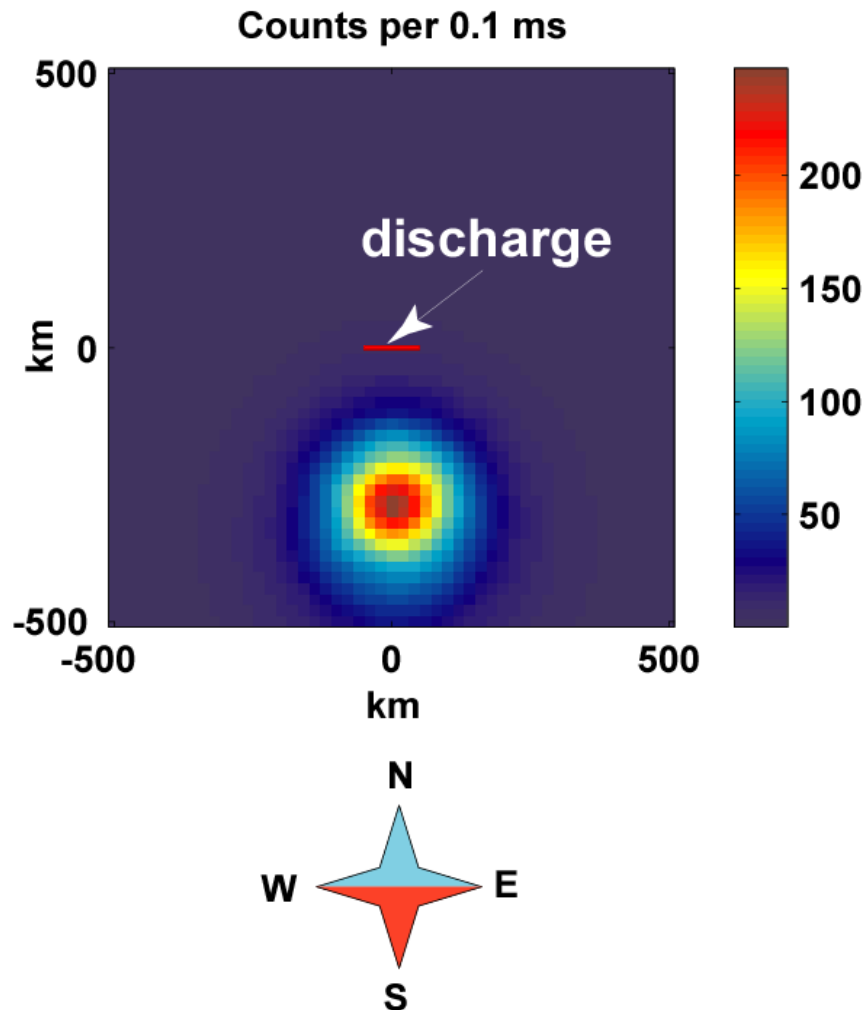


Runaway density:
time evolution





Simulated BATSE photon counts



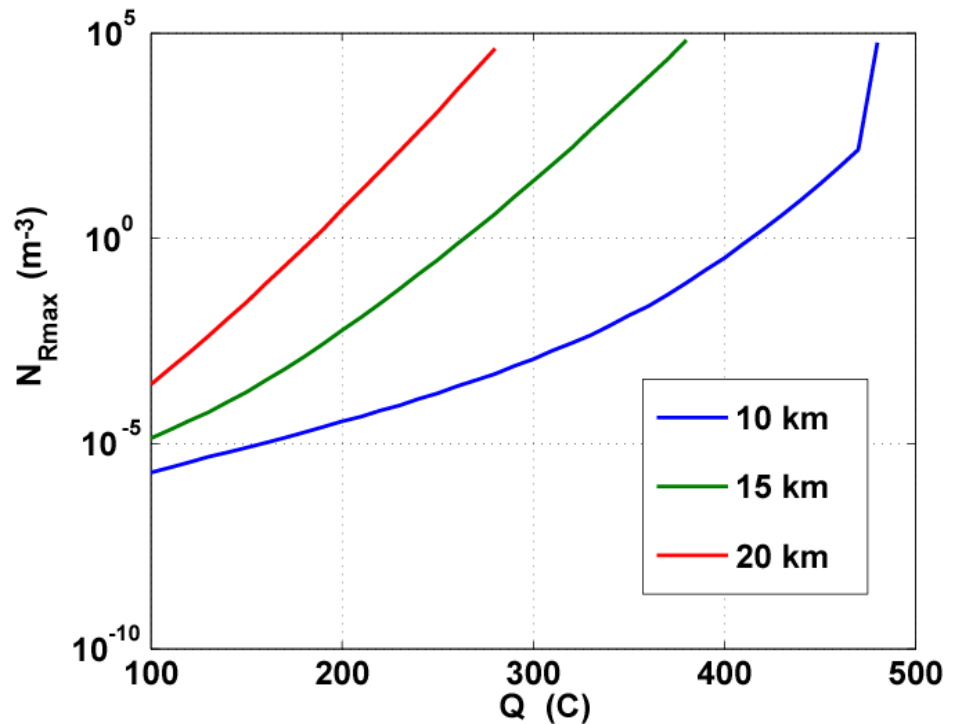
- $Q=1200$ C, 100 km long cloud (in EW direction)
- 45° N latitude
- 500 km altitude
- 100-300 keV
- Beam width $\sim mc^2/E_{el}$
- Spectrum is E_{ph}^{-1} on average, harder in the center, softer on the edges



Highly nonlinear dependence on the charge removed and altitude

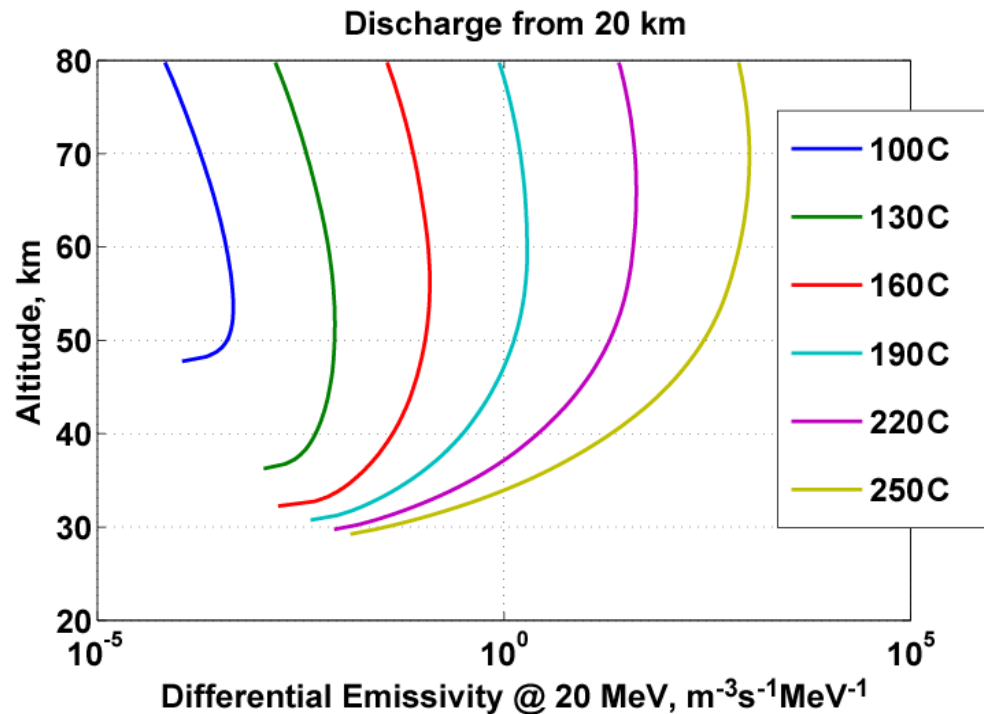
$$v \frac{dN}{dz} = RN + S$$

- $v=0.9c$ – runaway drift velocity
- $R(E/E_t) \sim N_m$ – avalanche rate
- $S \sim N_m$ – source from cosmic rays, $=10 \text{ m}^{-3} \text{ s}^{-1}$ at 10 km





Gamma Ray Emissivity



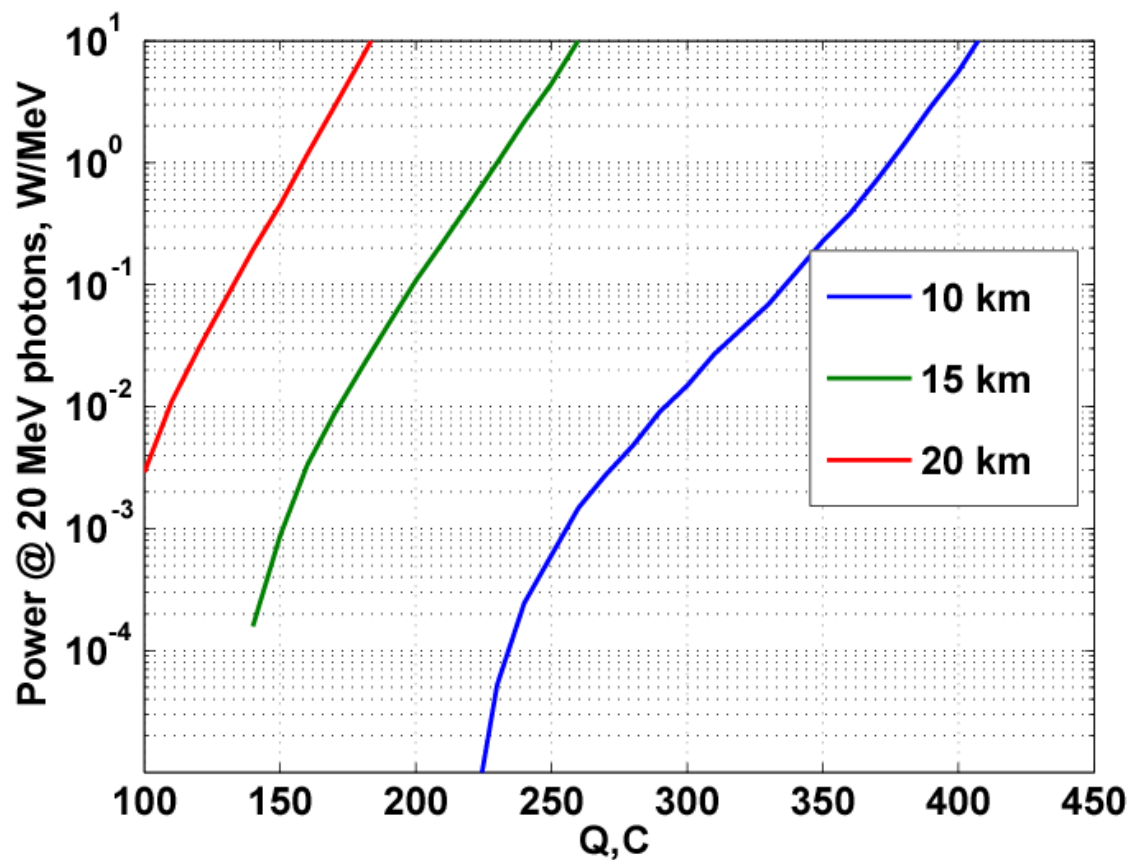
- $\varepsilon = N N_m v d\sigma/dE_{ph}$
- $d\sigma/dE_{ph}$ – Heitler's differential bremsstrahlung cross-section





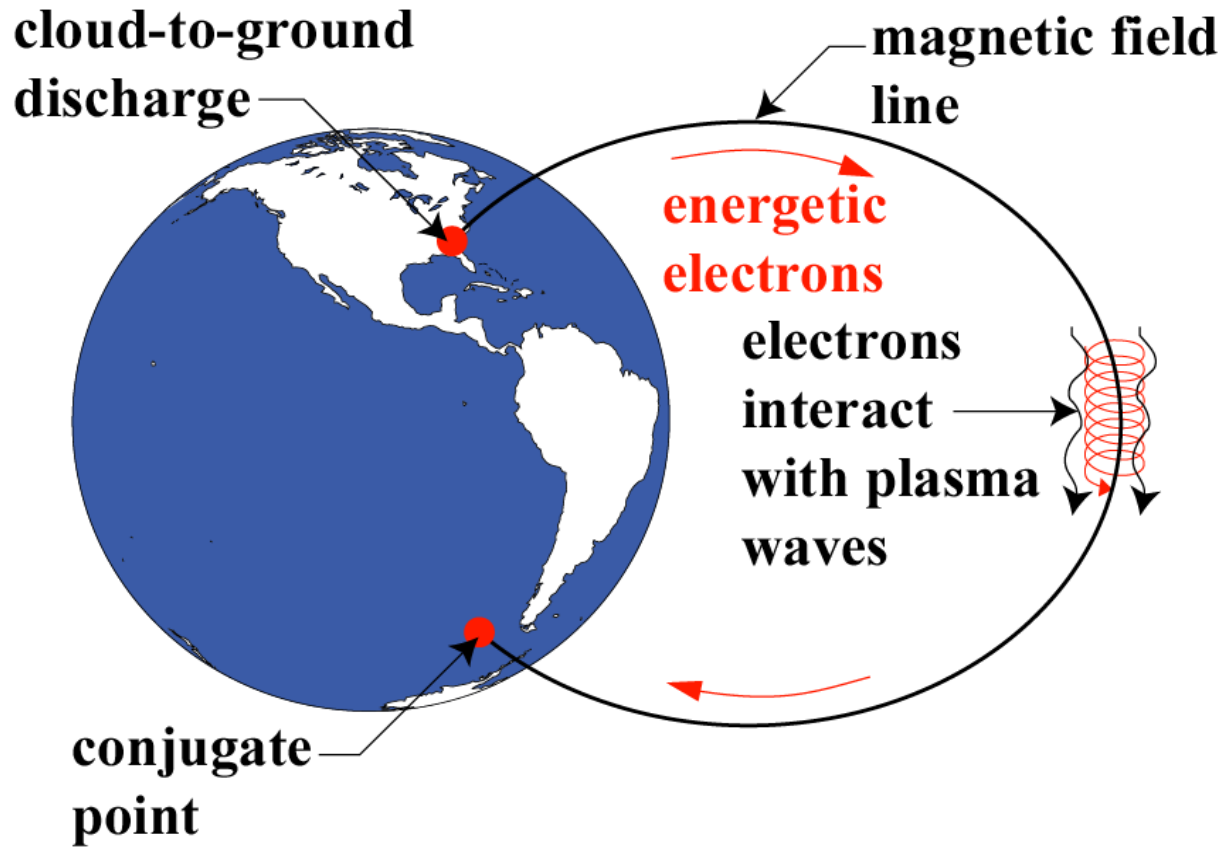
The power produced in gamma rays

Integrate over volume, assume $\sim 10\text{km}$ transverse size of the electron beam





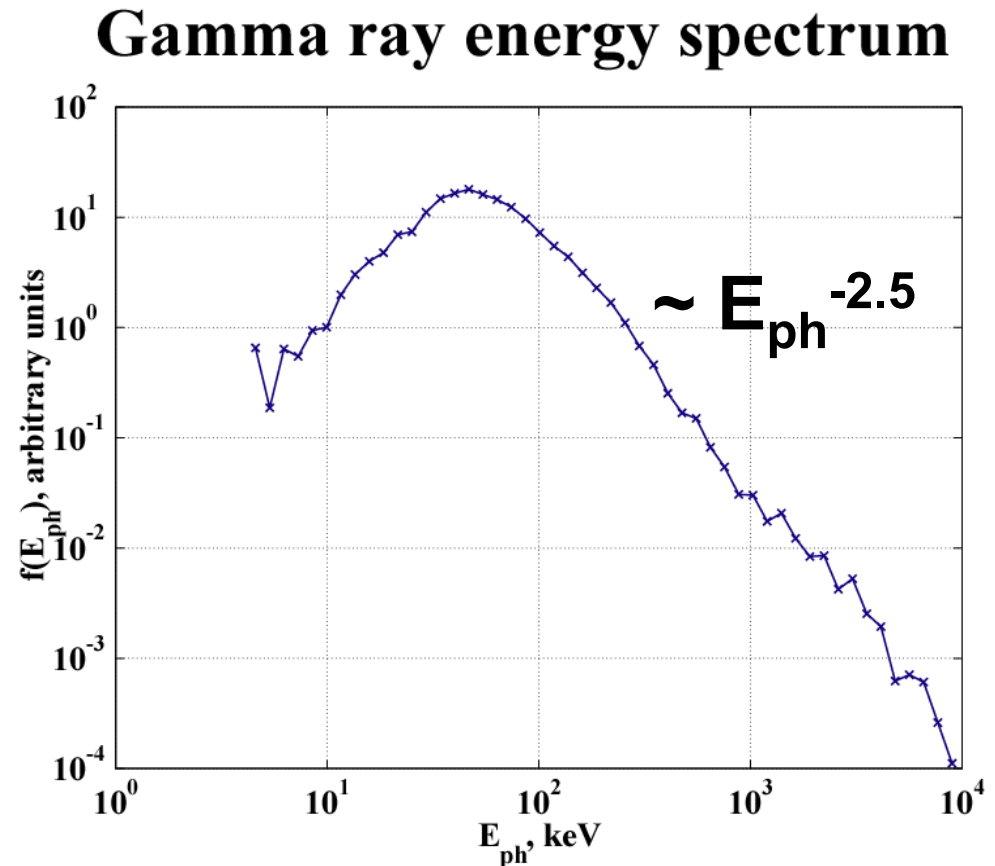
Escaping runaway electrons precipitate at the conjugate point





Gamma rays from precipitating electrons

- Initial gamma ray direction is downward
- They are backscattered
- Soft spectrum
- Observed spectrum is $\sim E_{ph}^{-1}$





Runaway Electron Avalanche and TGFs at the Magnetic Equator



Relativistic Runaway Breakdown at the Equator

- Numerous RHESSI observations

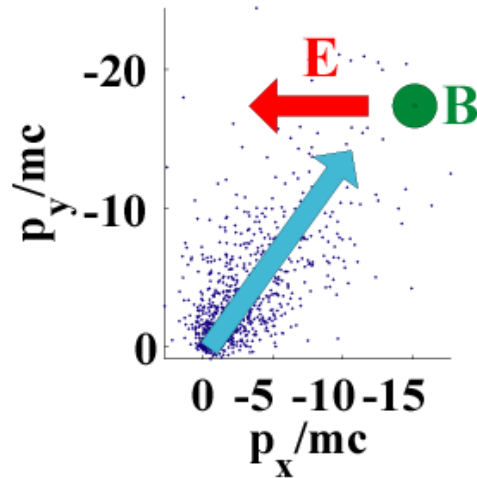
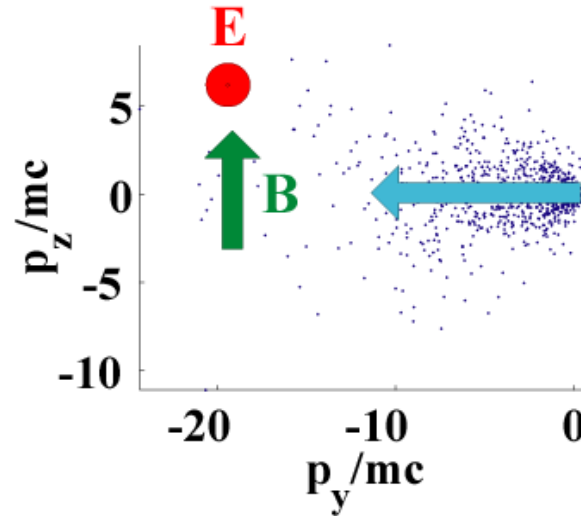
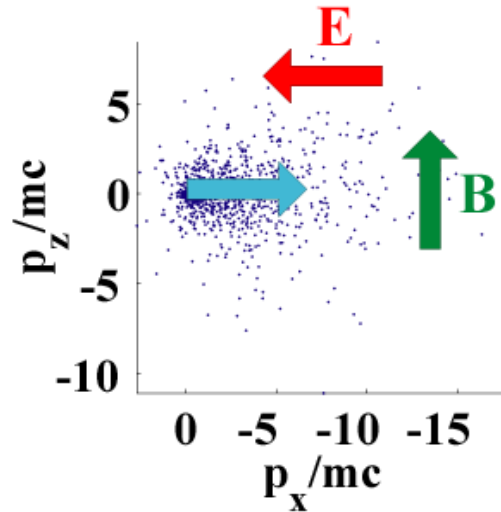
Contradiction:

- Suppressed by geomagnetic field
- \mathbf{ExB} drift in the horizontal direction
=> gamma rays are not emitted
vertically?



Electrons in momentum space

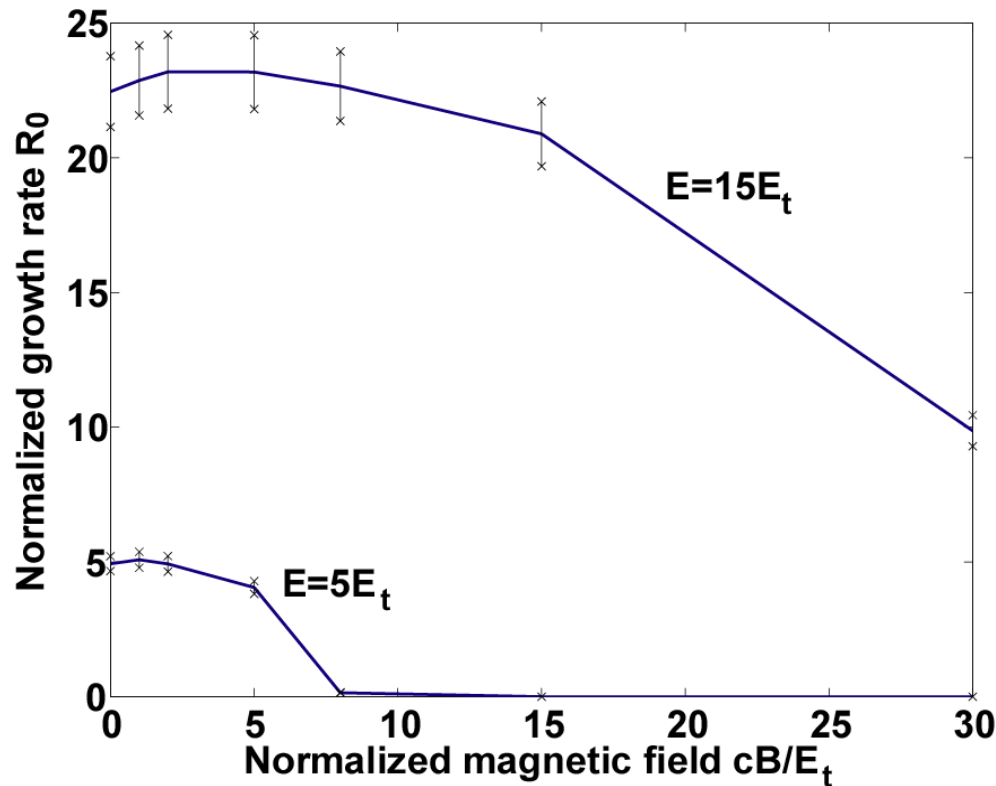
$E = 5 E_t$, $B = cE$, $B \perp E$



NOTE: There is no avalanche for $B \perp E$, $B=2cE$



The avalanche growth rates with perpendicular magnetic field



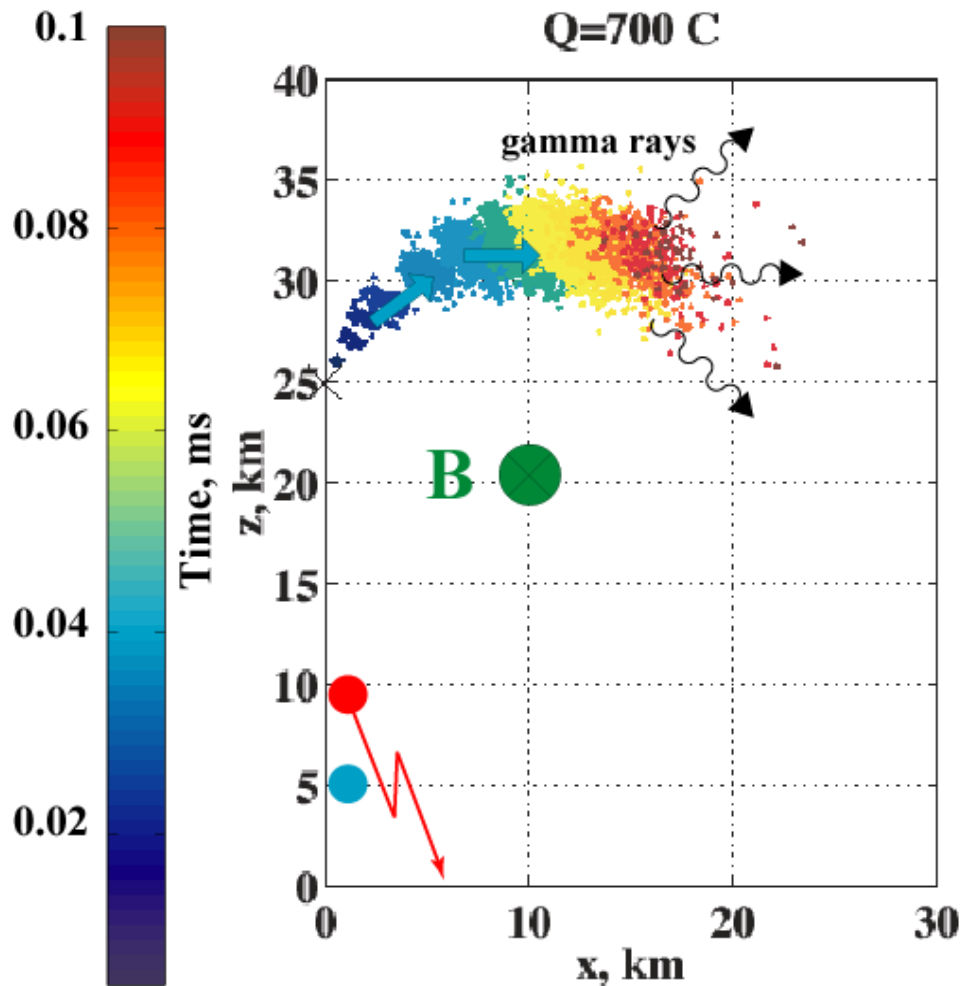
- The rate decreases slowly with increasing B at small B
- The avalanche is quenched approximately when $B > 2E/c$



Nonuniform Monte Carlo results at the equator

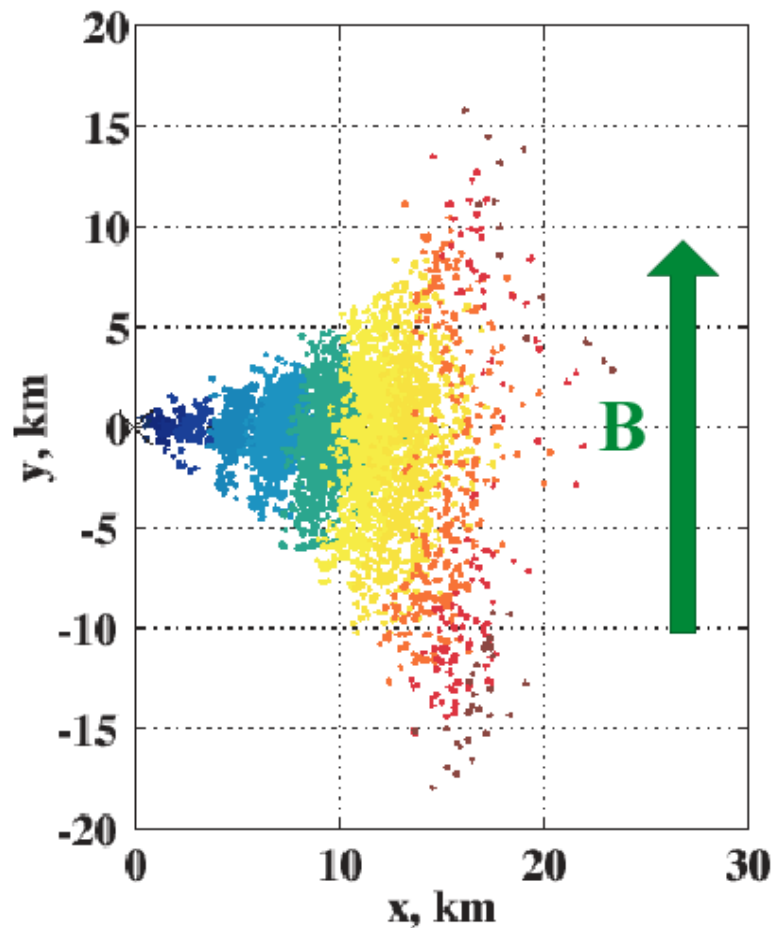
Side view

$Q=700\text{ C}$



Top view

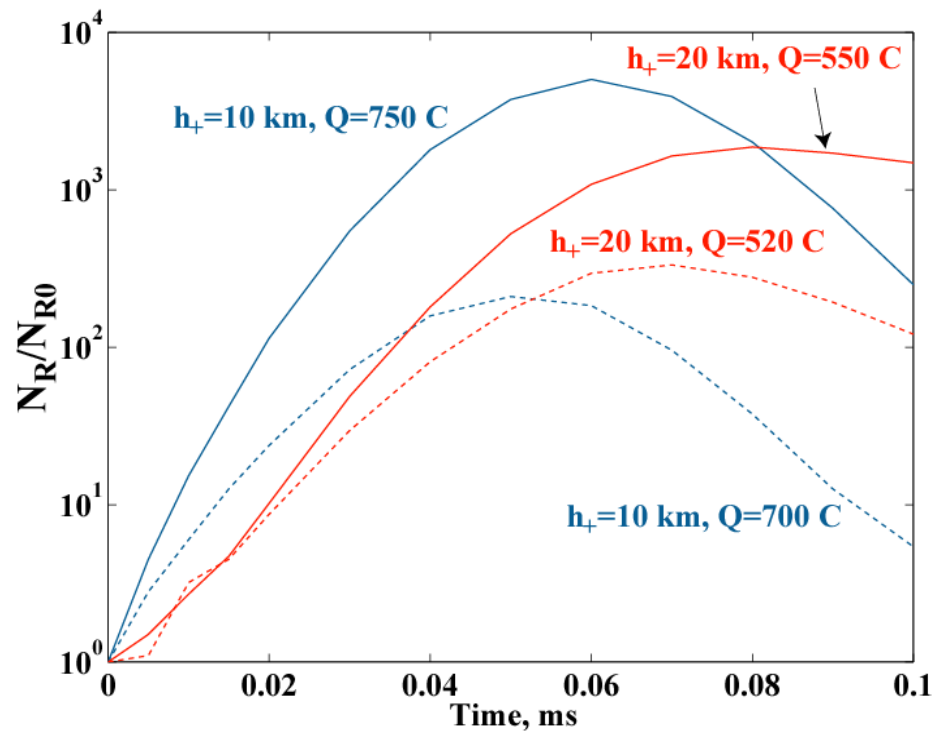
$Q=700\text{ C}$





Number of particles (MC simulation at the equator)


Number of particles with energy > 2 keV

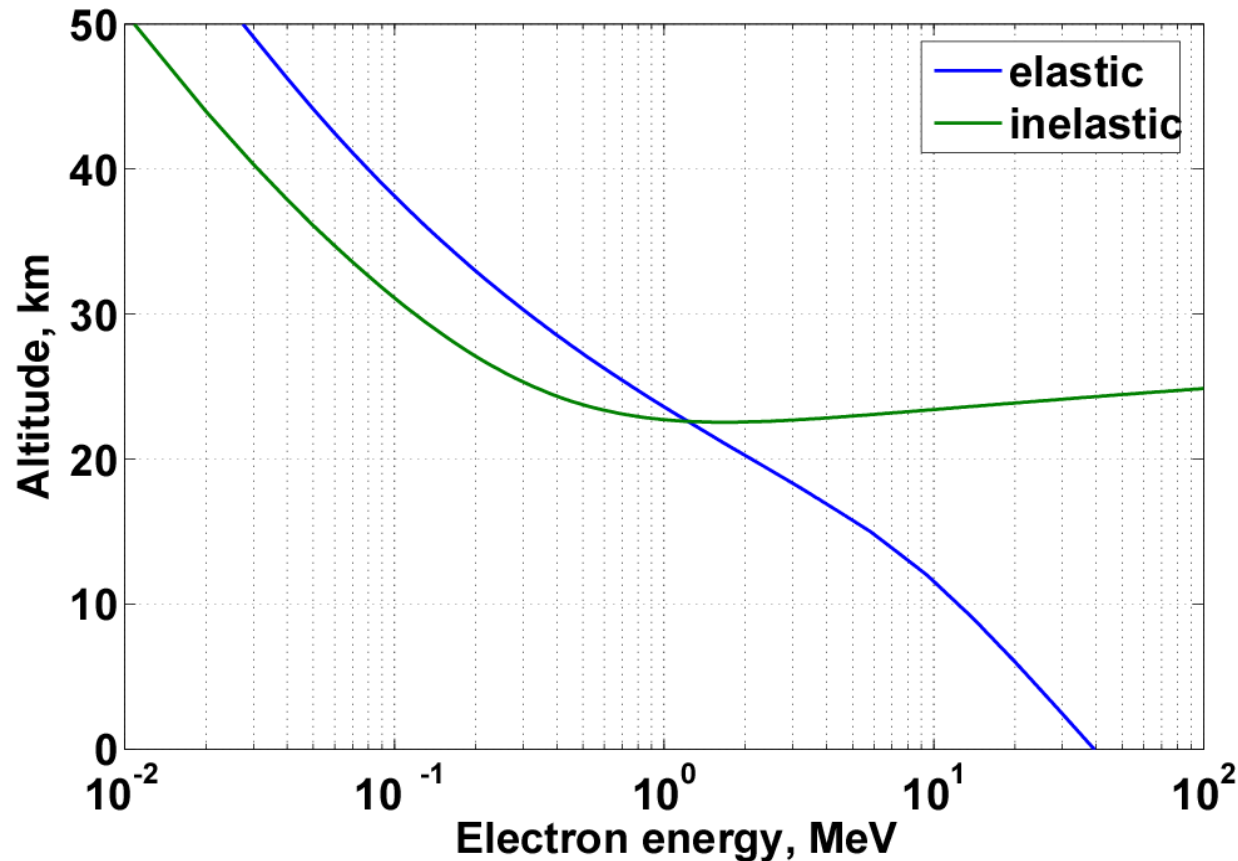


- Higher charge removal needed
- Energetic electrons are lost quickly due to moving out of the high E field region



Altitude above which B becomes important

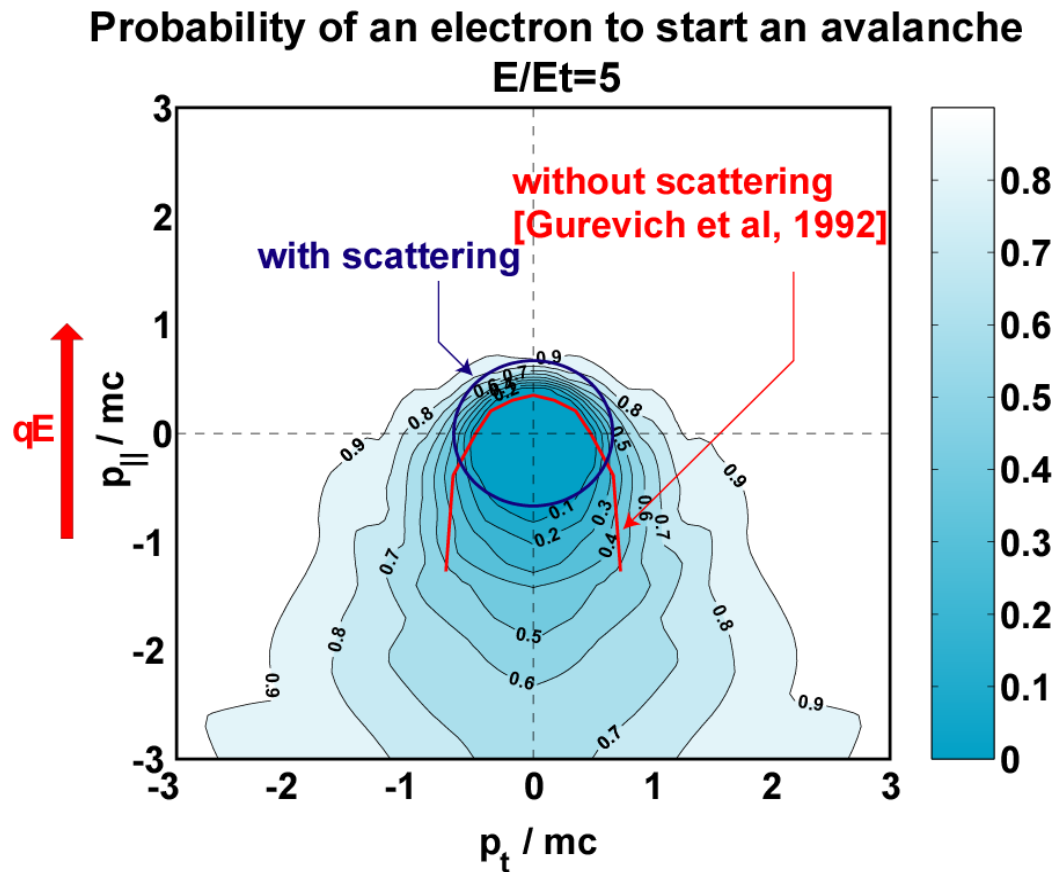
- For $f_{H0}=1$ MHz
- Elastic: $\nu_m = \omega_H$ (ν_m is the momentum transfer rate) 
- Inelastic: $F_D/p = \omega_H$ (F_D is the dynamic friction)





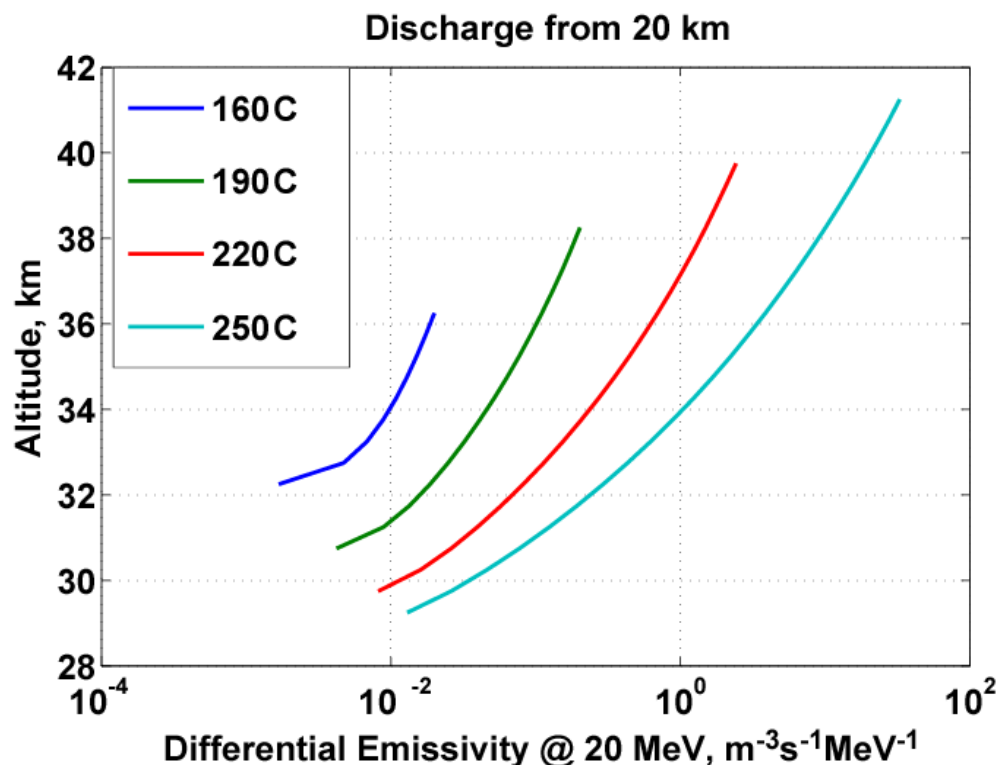
Energy of avalanche-producing electrons

- Low energy electrons are stopped by friction
- Avalanche continues only if $E > E_{\min}$
- $E_{\min} \sim mc^2/R_0$
- For a uniform avalanche, there are more low-energy electrons $\Rightarrow E_{\min}$ is the important energy scale





Gamma Ray Emissivity at the Equator

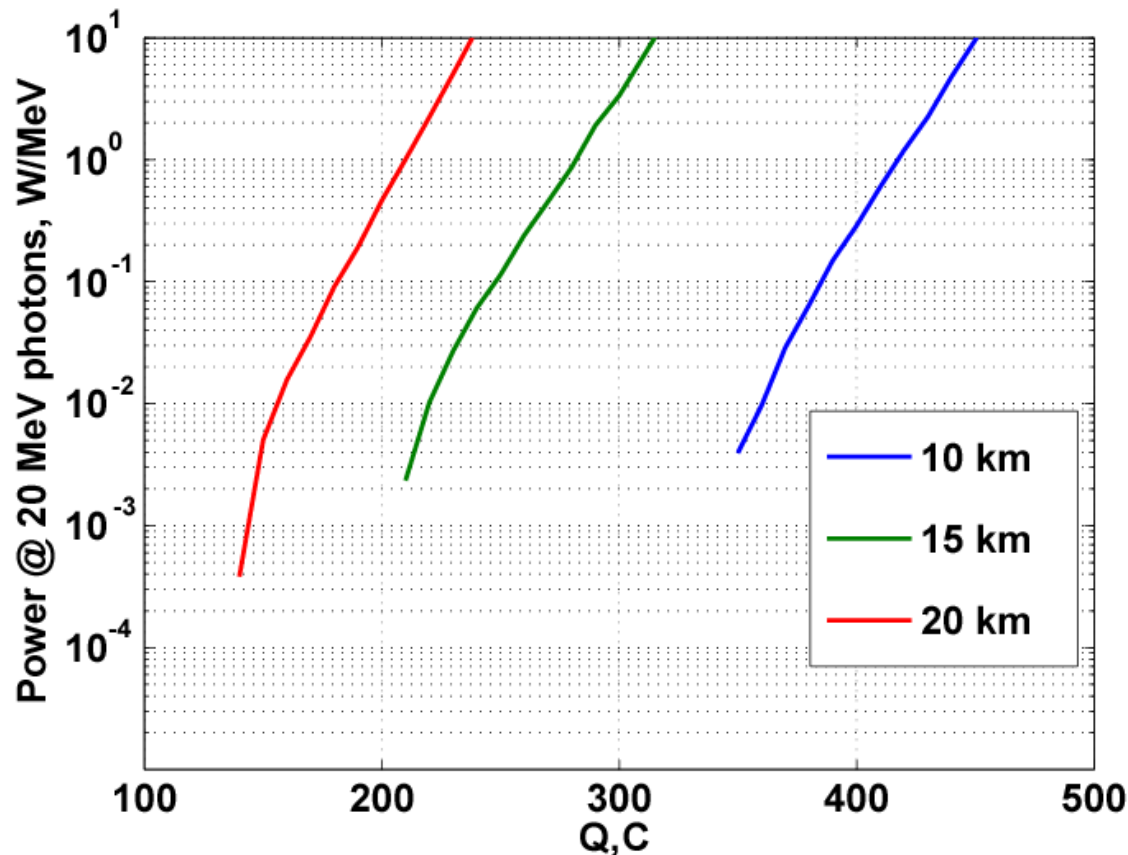


- $\varepsilon = N N_m v d\sigma/dE_{ph}$
- $d\sigma/dE_{ph}$ – Heitler's bremsstrahlung cross-section
- Avalanche is suppressed at high altitudes (when $B > 2E/c$)



The power produced in gamma rays at the equator

Integrate over volume, assume $\sim 10\text{km}$ transverse size of the electron beam





Conclusions

- **The number of relativistic electrons produced in the avalanche depends very nonlinearly on the electric field**
- **At the equator, the avalanche has to take place below 40 km** 